Dimensional analysis and Similarity

Problems and solutions

Laboratory experiment Drag force Similarity $\mathcal V$ **Prototype** \bar{U} $\delta(x)$ $L_p = 40 \; m$ $\boldsymbol{\mathcal{X}}$ **Similarity Model ?** $C_D = f(R_L)$ $\delta(x)$ $L_m = 2 m$

Reynold similarity;
$$
R_L|_p = R_L|_m = R_L
$$
 $C_D|_p = C_D|_m$

Same fluid $\mu_m = \mu_p = \mu$, and $\rho_m = \rho_p = \rho$

Different fluids $\mu_m \neq \mu_p$, and $\rho_m \neq \rho_p$

Worked example 11

If the rotating cone has a base diameter of 120mm and is to be rotated first in ambient air and subsequently in water. What rotational speed should be used in water to obtain the same Reynolds number as that obtained in air at 3000 rpm? Take the density and viscosity of air to be 1.2 kg/m³ and $1.8x10^{-5}$ kg/ms respectively and the density and viscosity of water to be 1000 kg/m³ and 0.001 kg/ms respectively.

Answer: 200 rpm

Reynolds similarity $Re = \frac{\rho \omega D^2}{\mu}$

$$
Re_{air} = \frac{1.2 \times (3000 \times \frac{2\pi}{60}) \times 0.12^2}{1.8 \times 10^{-5}} = 3.02 \times 10^5
$$

$$
Re_{water} = \frac{1000 \times (N \times \frac{2\pi}{60}) \times 0.12^2}{0.001}
$$

= 3.02 × 10⁵

Relation between air and water rotational speed

$$
Re_{air} = \left(\frac{\rho \omega D^2}{\mu}\right)_{air} = Re_{water} = \left(\frac{\rho \omega D^2}{\mu}\right)_{water}
$$

$$
Re_{air} = \left(\frac{\rho N \left\{\frac{2\pi}{60}\right\} D^2}{\mu}\right)_{air} = Re_{water} = \left(\frac{\rho N \left\{\frac{2\pi}{60}\right\} D^2}{\mu}\right)_{water} \qquad \left(\frac{\rho N}{\mu}\right)_{air} = \left(\frac{\rho N}{\mu}\right)_{water}
$$

$$
N_{water} = N_{air} \times \frac{\rho_{air}}{\rho_{water}} \frac{\mu_{water}}{\mu_{air}} = \frac{N_{air}}{15}
$$

Use to build our model

In the present example

 N_{air} = 3000 rpm then N_{water} = N = 200 rpm

Worked example 12

The torque on the rotating cone shown is related to the rotation speed, fluid properties and cone base diameter:

 $T = f(\omega, \rho, \mu, D)$

Identify the appropriate non-dimensional groups.

Answer: Cm, Re

 $\overrightarrow{T} = f(\omega, \rho, \mu, D)$

 \bullet First find the dimensions of each variable

- 5 variables (n = 5)
- 3 dimensions $(k = 3)$
- M=5-3=2 dimensionless groups
- I choose my two non repeating variables $\left(\!\frac{1}{2}\right)$ and μ
- The repeating variables used to form the Π -groups are ρ , D and ω . These cannot be used to form a non-dimensional group on their own and so are suitable.

Choose T because it is the output variable

 $Re = \frac{\rho \omega D}{\sqrt{v}}$

 \mathbf{z}

• The two Π groups are

$$
[\Pi_1] = [\mathbf{T}][\rho]^a[\mathbf{D}]^b[\omega]^c = [\mathbf{M}^0 \mathbf{L}^0 \mathbf{T}^0]
$$

$$
[\Pi_1] = [\mathbf{M} \mathbf{L}^2 \mathbf{T}^{-2}][\mathbf{M} \mathbf{L}^{-3}]^a[\mathbf{L}]^b[\mathbf{T}^{-1}]^c = [\mathbf{M}^0 \mathbf{L}^0 \mathbf{T}^0]
$$

• Remember maths from way back – when you multiply you add the exponents?

$$
x^ax^b=x^{a+b}
$$

- We can now look at each of the dimensions in turn and balance the exponents
- For example taking [M]
- $1 + a = 0$ so $a = -1$

$$
M^1M^a=M^0
$$

 \overline{a}

- Similarly for L and T
- L: $2-3a+b=0$
- T: $-2 c = 0$ so c=-2
- Thus $b = -5$ (from the L equation)
- We have $[\Pi_1] = [T][\rho]^a[D]^b[\omega]^c$

• So
$$
\Pi_1 = T \rho^{-1} D^{-5} \omega^{-2} = \frac{T}{\rho \omega^2 D^5}
$$

- Compare this to the torque coefficient
- Buckingham- Π cannot identify constants.
- Our first non-dimensional group is Cm

$$
Cm = \frac{T}{\frac{1}{2}\rho\omega^2 D^5}
$$

$$
C_m = \frac{T}{\frac{1}{2}\rho(\omega D)^2 D^2 D}
$$

• We can follow the same process for the second non-dimensional group:

$$
[\Pi_2] = [\mu][\rho]^d[D]^e[w]^f = [M^0L^0T^0]
$$

$$
[\Pi_2] = [ML^{-1}T^{-1}][ML^{-3}]^d[L]^e[T^{-1}]^f = [M^0L^0T^0]
$$

- Again balancing the exponents:
- $M: 1+d=0$ so d=-1
- $L: -1-3d+e=0$
- $T: -1-f=0$ so $f=-1$
- Thus e=-2

$$
\Pi_2 = \mu \rho^{-1} D^{-2} \omega^{-1} = \frac{\mu}{\rho \omega D^2}
$$

As we saw before *Re* **is one of the dominant similar parameters in the governing equations and need to be included in the analysis**

- Compare this to the rotating Reynolds number
- Π_2 is 1/Re (which is also dimensionless of course).

 $Re = \frac{\rho \omega D^2}{2}$

The functional relationship is therefore *Cm=f(Re)*

Problem:

On worked example 12 we are looking for the measurement of the **pressure distribution** inside the cone, where the pressure on the rotating cone is related to the rotation speed, fluid properties, viscosity μ and density ρ , and cone base diameter:

$$
P=f(\omega,\rho,\mu,D)
$$

- 5 variables ($n = 5$)
- 3 dimensions $(k = 3)$
- M=5-3=2 dimensionless groups
- I choose my two non repeating variables as P (objective function) and μ
- The repeating variables used to form the Π -groups are ρ , D and ω .

$$
\pi_1 = P \rho^{a1} \omega^{a2} D^{a3}
$$

\n
$$
\pi_1: [ML^{-1}T^{-2}][ML^{-3}]^{a1}[T^{-1}]^{a2}[L]^{a3} = [M^0L^0T^0]
$$

\n
$$
M: 1 + a1 = 0; a1 = -1
$$

\n
$$
L: -1 - 3a1 + a3 = 0; 2 + a3 = 0; a3 = -2
$$

\n
$$
T: -2 - a2 = 0; a2 = -2
$$

\n
$$
\pi_1 = P \rho^{-1} \omega^{-2} D^{-2} = \frac{P}{\rho \omega^2 D^2} = E_u
$$

$$
\pi_2 = \mu \rho^{b1} \omega^{b2} D^{b3}
$$

\n
$$
\pi_2: [ML^{-1}T^{-1}][ML^{-3}]^{b1}[T^{-1}]^{b2}[L]^{b3} = [M^0L^0T^0]
$$

\n
$$
M: 1 + b1 = 0; b1 = -1
$$

\n
$$
L: -1 - 3b1 + b3 = 0; 2 + b3 = 0; b3 = -2
$$

\n
$$
T: -1 - b2 = 0; b2 = -1
$$

\n
$$
\pi_2 = \mu \rho^{-1} \omega^{-1} D^{-2} = \frac{\mu}{\rho \omega D^2} = \frac{1}{R_e}
$$

Therefore: $E_u = f(R_e)$

Problem:

On worked example 12 the working fluid is air and the **rotational velocity is large enough to compressibility been significant**.

As before we are looking the **torque** on the cone; the torque on the rotating cone is related to the rotation speed, fluid properties, viscosity μ and density ρ , cone base diameter, and this time as the fluid is compressible we also need to consider the pressure P and take into account that the speed of sound is $c = \sqrt{\frac{kp}{n}}$ ρ , i.e. $p=\frac{1}{\mu}$ κ ².

$$
\boldsymbol{T} = f(\omega, \rho, \boldsymbol{\mu}, D, \boldsymbol{P})
$$

Thus we expect $n - k = 6 - 3 = 3$ Pi groups:

$$
\pi_1 = D \rho^{a_1} \omega^{a_2} D^{a_3}
$$
 From example 12:
$$
Cm = \frac{T}{\frac{1}{2} \rho \omega^2 D^5}
$$
Thermofluids 2

$$
\pi_2 = \mu \rho^{b1} \omega^{b2} D^{b3}
$$
 From example 12: $\pi_2 = \frac{\mu}{\rho \omega D^2} = \frac{1}{R_e}$

$$
a = P \rho^{c1} \omega^{c2} D^{c3}
$$
 From the previous example: $\pi_3 = \frac{P}{\rho \omega^2 D^2}$

Taking into account the **speed of sound** is $c = \sqrt{\frac{kp}{q}}$ ρ , i.e**.** $\boldsymbol{p}=\frac{1}{L}$ \boldsymbol{k} \mathbf{z}

$$
\pi_3 = \frac{c^2}{k\omega^2 D^2} = \frac{1}{kM_a^2};
$$
 compressible fluid $E_u = \frac{1}{kM_a^2}$

$$
C_m = f(R_e, \sqrt{k} M_a
$$

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 $\pi_3 = E_u$

Problem:

On a **Boundary layer** over a flat plate, the **drag force** D varies with the characteristic length of the plate L, freestream velocity U, viscosity μ , and density ρ . Find the dimensionless parameters for this problem.

Solution: The functional relationship is $\mathbf{D} = f(U, \mu, \rho, L)$, with $n = 5$ variables and $k = 3$ primary dimensions (M,L,T).

Thus we expect $n - k = 5 - 3 = 2$ Pi groups:

$$
\pi_1 = D \rho^{a1} U^{a2} L^{a3}
$$

$$
\pi_1: [MLT^{-2}][ML^{-3}]^{a1}[LT^{-1}]^{a2}[L]^{a3} = [M^{0}L^{0}T^{0}]
$$

\nM: 1 + a1 = 0; a1 = -1
\nL: 1 - 3a1 + a2 + a3 = 0; 4 + a2 + a3 = 0
\nT: -2 - a2 = 0; a2 = -2; a3 = -2
\n
$$
\pi_1 = D \rho^{-1}U^{-2}L^{-2} = \frac{D}{\rho U^2 L^2} = \frac{D}{\rho U^2 A} = C_D/2
$$

$$
\pi_2 = \mu \rho^{b1} U^{b2} L^{b3}
$$

\n
$$
\pi_2: [ML^{-1}T^{-1}][ML^{-3}]^{b1}[LT^{-1}]^{b2}[L]^{b3} = [M^{0}L^{0}T^{0}]
$$

\n
$$
M: 1 + b1 = 0; b1 = -1
$$

\n
$$
L: -1 - 3b1 + b2 + b3 = 0; 2 + b2 + b3 = 0
$$

\n
$$
T: -1 - b2 = 0; b2 = -1; a3 = -1
$$

\n
$$
\pi_2 = \mu \rho^{-1} U^{-1} L^{-1} = \frac{\mu}{\rho U L} = \frac{1}{R_L}
$$

Therefore:

$$
C_D = f(R_L)
$$

Thus $C_D = f(R_L)$ just as in the boundary layer theory.

As can be observed from the above analysis, this result is also valid for any type of boundary layer.

Problem:

In an atomic explosion, a rapid release of a significant amount of energy occurs within a small region. A strong spherical shock wave develops at the point of detonation: in the early state, the pressure behind the wave front is several thousand times the initial air pressure, whose influence may neglected. Thus the radius of the shock wave front, R , at an interval of time, T, after the explosion depends on the quantity of energy release, E, the time, T, and on the initial air density, ρ .

Photograph of the fireball of the atomic explosion in New Mexico

Taylor (1941) analysed data from a series of high-speed photographs of the expansion of a fireball taken during an American nuclear test and using the π theorem concluded that the radius of the wave front satisfies the following relation:

$$
R = \pi E^{1/5} T^{2/5} \rho^{-1/5}
$$

The dimensions of the above quantities are:

$$
R[L]; T[T]; \rho[M L^{-3}]; E[ML^{2}T^{-2}]
$$

Using the π theorem repeat Taylor's analysis.

$$
m = n - k = 4 - 3 = 1
$$

\n
$$
\pi = RT^a \rho^b E^c \rightarrow [L][T]^a [ML^{-3}]^b [ML^2 T^{-2}]^c = [L^0 T^0 M^0]
$$

\n
$$
L: 1 - 3b + 2c = 0
$$

\n
$$
T: a - 2c = 0
$$

\n
$$
M: b + c = 0
$$

\nSolution: $a = -\frac{2}{5}, b = \frac{1}{5}, c = -\frac{1}{5}$, therefore:
\n
$$
\pi = RT^{-2/5} \rho^{1/5} E^{-1/5}, \qquad or \quad \mathbf{R} = \pi \frac{T^{2/5} E^{1/5}}{\rho^{1/5}}
$$

with π as experimental constant.

The power input to a centrifugal pump is a function of the volume flowrate, Q, impeller diameter D, rotation rate Ω and fluid properties μ and ρ .

$$
P = f(Q, D, \Omega, \rho, \mu)
$$

Objective of the experiment

To determine P as a function Q by a *Re* similarity

Re-write this as a dimensionless relationship using Ω , ρ and D as the repeating variables.

$$
\frac{P}{\rho \Omega^3 D^5} = f\left(\frac{\mu}{\rho \Omega D^2}, \frac{Q}{\Omega D^3}\right)
$$

 $P = f(Q, D, \Omega, \rho, \mu)$

• First find the dimensions of each variable

- 6 variables $(n = 6)$
- 3 dimensions $(k = 3)$
- m=6-3=3 dimensionless groups
- I choose my three variables as P, Q and μ
- The variables used to form the Π -groups are ρ , D and Ω . These cannot be used to form a non-dimensional group so are suitable.

 $[\Pi_1] = [P][\rho]^a[D]^b[\Omega]^c = [M^0L^0T^0]$ $[\Pi_1] = [ML^2T^{-3}][ML^{-3}]^a[L]^b[T^{-1}]^c = [M^0L^0T^0]$

- $M: 1+a=0$ so $a = -1$
- L: $2-3a+b=0$
- T: -3 -c=0 so c= -3
- Thus $b = -5$ (from the L equation)

$$
\Pi_1 = P \rho^{-1} D^{-5} \Omega^{-3} = \frac{P}{\rho \Omega^3 D^5} = \frac{P}{\frac{1}{2} \rho (\Omega D)^2 D^2 (\Omega D)}
$$

• Likewise for Π_2

 $[\Pi_2] = [\mu][\rho]^d[D]^{e}[\Omega]^{f} = [M^0L^0T^0]$ $[\Pi_2] = [ML^{-1}T^{-1}][ML^{-3}]^d[L]^e[T^{-1}]^f = [M^0L^0T^0]$

- $M: 1+d=0$ so $d=-1$
- $L: -1-3d+e=0$
- T: $-1-f=0$ so $f=-1$
- \bullet Thus e=-2

$$
\Pi_2 = \mu \rho^{-1} D^{-2} \Omega^{-1} = \frac{\mu}{\rho \Omega D^2}
$$

• Again this is 1/Re, the rotating Reynolds number

• And for Π_3

$$
[\Pi_3] = [Q][\rho]^g[D]^h[\Omega]^i = [M^0L^0T^0]
$$

$$
[\Pi_3] = [L^3T^{-1}][ML^{-3}]^g[L]^h[T^{-1}]^i = [M^0L^0T^0]
$$

- \bullet M: $g=0$
- L: 3-3g+h=0 and as g=0 then h=-3
- T: -1 -i=0 so i= -1

$$
\Pi_3 = \mathbf{Q} \rho^0 \mathbf{D}^{-3} \mathbf{\Omega}^{-1} = \frac{\mathbf{Q}}{\mathbf{\Omega} \mathbf{D}^3}
$$

• Thus

$$
\frac{P}{\rho \Omega^3 D^5} = f\left(\frac{\mu}{\rho \Omega D^2}, \frac{Q}{\Omega D^3}\right)
$$

• As will be seen in the Turbomachinery topic, these are the standard nondimensional groups used for correlating pump power.

]

What about dimensionless variables?

- Sometimes there is a variable that affects the outcome but has no dimensions
	- Eg the number of blades on a wind turbine
	- Or the angle of attack on an aerofoil
- This is added in as another variable but is automatically one of the non-dimensional groups
- Looking at WE13 for example, add number of blades as a variable
	- Number of variables increases from 6 to 7
	- Number of dimensions remains 3
	- Number of groups increases from 3 to 4
	- Additional group is N
	- Functional relationship becomes:

$$
\frac{P}{\rho \Omega^3 D^5} = f\left(\frac{\mu}{\rho \Omega D^2}, \frac{Q}{\Omega D^3}, N\right)
$$

Problem:

A liquid of density ρ and viscosity μ flows in a channel of width H with a characteristic velocity U, one small side of the upper wall is heated at a temperature T_w with respect to the inlet flow temperature T_0 while the rest of the walls are thermally isolated. Knowing that thermal conductivity of the liquid k has dimension $\left[M\ L\ T^{-3}\Theta^{-1} \right]$ and its specific heat c_p has dimension $\left[L^2|T^{-2}\Theta^{-1}\right]$ find a dimensionless relation for the wall **heat flux** q of dimension $[M T^{-3}].$

$$
\mathbf{q} = f(\mathbf{k}, c_p, \rho, \boldsymbol{\mu}, H, U, (\mathbf{T_w} - \mathbf{T_0}))
$$

Units (dimensions) $\rightarrow MLT$ Θ ,

$$
m = n - k = 8 - 4 = 4
$$

Quantity of interest (q) and free parameters (k , μ , $(T_w - T_0)$) in the characteristic dimensionless numbers **(Péclet, Reynolds and dimensionless heat radiation).**

We are looking for a dimensionless heat flux in terms of the imposed dimensionless increase of temperature, and according to the governing equations we expect similarity with respect to the Reynolds and Peclet numbers.

$$
\pi_1 = q \rho^{a1} U^{a2} H^{a3} c_p^{a4}
$$

\n
$$
\pi_1: [M T^{-3}][ML^{-3}]^{a1}[LT^{-1}]^{a2}[L]^{a3}[L^2 T^{-2} \Theta^{-1}]^{a4} = [M^0 L^0 T^0]
$$

\n
$$
M: 1 + a1 = 0; \quad a1 = -1
$$

\n
$$
L: 0 - 3a1 + a2 + a3 + 2a4 = 0; \quad 3 + a2 + a3 + 2a4 = 0; \quad a3 = 0
$$

\n
$$
T: -3 - a2 - 2a4 = 0; \quad a2 = -3 \text{ T}
$$

\n
$$
\Theta: 0 - a4 = 0; \quad a4 = 0 \text{ T}
$$

\n
$$
\pi_1 = q \rho^{-1} U^{-3} = \frac{q}{\rho U^3}
$$

$$
\pi_2 = k \rho^{b1} U^{b2} H^{b3} c_p^{b4}
$$

\n
$$
\pi_2: [MLT^{-3}\Theta^{-1}] [ML^{-3}]^{b1} [LT^{-1}]^{b2} [L]^{b3} [L^2 T^{-2}\Theta^{-1}]^{b4} = [M^0 L^0 T^0]
$$

\n
$$
M: 1 + b1 = 0; b1 = -1
$$

\n
$$
L: 1 - 3b1 + b2 + b3 + 2b4 = 0; 4 + b2 + b3 + 2b4 = 0; b3 = -1
$$

\n
$$
T: -3 - b2 - 2b4 = 0; b2 = -1 \text{ T}
$$

\n
$$
\Theta: -1 - b4 = 0; b4 = -1 \text{ T}
$$

\n
$$
\pi_2 = k \rho^{-1} U^{-1} H^{-1} c_p^{-1} = \frac{k}{\rho U H c_p}; D = \frac{k}{\rho c_p} \rightarrow \pi_2 = \frac{D}{U H} = \frac{1}{P_e}
$$

$$
\pi_3 = \mu \rho^{c1} U^{c2} H^{c3} c_p^{c4}
$$
\n
$$
\pi_3: [M L^{-1} T^{-1}] [ML^{-3}]^{c1} [LT^{-1}]^{c2} [L]^{c3} [L^2 T^{-2} \Theta^{-1}]^{c4} = [M^0 L^0 T^0]
$$
\n
$$
M: 1 + c1 = 0; c1 = -1
$$
\n
$$
L: -1 - 3c1 + c2 + c3 + 2c4 = 0; 2 + c2 + c3 + 2c4 = 0; c3 = -1
$$
\n
$$
T: -1 - c2 - 2c4 = 0; c2 = -1 \quad 1
$$
\n
$$
\Theta: 0 - c4 = 0; c4 = 0 \quad 1
$$
\n
$$
\pi_3 = \mu \rho^{-1} U^{-1} H^{-1} = \frac{\mu}{\rho U H}; \pi_3 = \frac{1}{R_e}
$$
\n
$$
\dots
$$
\n
$$
\pi_4 = (T_w - T_0) \rho^{d1} U^{d2} H^{d3} c_p^{d4}
$$
\n
$$
\pi_2: [\Theta] [ML^{-3}]^{d1} [LT^{-1}]^{d2} [L]^{d3} [L^2 T^{-2} \Theta^{-1}]^{d4} = [M^0 L^0 T^0]
$$
\n
$$
M: 0 + d1 = 0; d1 = 0
$$
\n
$$
L: 0 - 3d1 + d2 + d3 + 2d4 = 0; d2 + d3 + 2d4 = 0; d3 = 0
$$
\n
$$
T: -d2 - 2d4 = 0; d2 = -2 \quad 1
$$
\n
$$
\Theta: 1 - d4 = 0; d4 = 1 \quad 1
$$
\n
$$
\pi_4 = (T_w - T_0) U^{-2} c_p = \frac{(T_w - T_0) c_p}{U^2}
$$
\nTherefore: $\frac{q}{\rho U^3} = f(R_e, P_e, \frac{(T_w - T_0) c_p}{U^2})$

Problem:

When tested in water at 20°C flowing at 2 m/s, an 8-cm-diameter sphere has a measured drag of 5 N ($kg~m~s^{-2}$). What will be the velocity and drag force on a 1.5**m-diameter weather balloon moored in sea-level standard air under dynamically similar conditions?**

Solution: For water at 20°C take $\rho = 998 \frac{kg}{m^3}$ $m_{\frac{1}{2}}^{3}$ and $\mu=0.001\frac{kg}{mc}$ $m s$. For sea-level standard air take $\rho = 1.2255 \frac{kg}{m^3}$ $m³$ and $\mu = 1.78 \ 10^{-5} \frac{kg}{m}$ $m s$.

The balloon velocity follows from *dynamic similarity*, which requires identical Reynolds numbers:

$$
R_e\Big|_m = R_e\Big|_p
$$

\n
$$
R_e\Big|_m = \frac{\rho UD}{\mu}\Big|_m = \frac{998x2x0.08}{0.001} = 1.6 10^5
$$

\n
$$
R_e\Big|_p = \frac{\rho UD}{\mu}\Big|_p = \frac{1.2255 x U_p x 1.5}{1.78 10^{-5}} = 1.6 10^5
$$

$$
or U_p = U_{\text{balloon}} = 1.55 \text{ m/s.}
$$

Then the two spheres will have identical drag coefficients:

$$
\left(\frac{C_D = f(R_e)}{\rho U^2 A} = 2C_D\right)
$$

$$
C_D\Big|_m = C_D\Big|_p
$$

$$
C_D\Big|_m = \frac{D}{\frac{1}{2}\rho U^2 D^2}\Big|_m = \frac{5}{\frac{1}{2}998 \times 2^2 \times 0.08^2} = 0.396
$$

$$
D\Big|_p = \frac{1}{2}\rho U^2 D^2\Big|_p C_D\Big|_m = \frac{1}{2} 1.2255 \times 1.55^2 \times 1.5^2 \times 0.396 = 1.3 N
$$

Therefore
$$
D_{\text{balloon}} = 1.3 \text{ N.}
$$
 The University of Nottingham

Worked example 14

A prototype boat is to be tested at a model scale of 50:1. This is an incompressible free surface flow and to ensure dynamic similarity both Reynolds number and Froude number must be equal. What kinematic viscosity must the model working fluid have if the prototype working fluid is water ($v=10^{-6}$ m²/s)?

Answer: $v_m = 2.83 \times 10^{-9} m^2/s$

- We want $Fr_m = Fr_p$ where $Fr = \frac{U^2}{gL}$ gL
- So $\frac{U_m^2}{g}$ gL_n U_p^2 $g L_{\bm p}$
- Geometric similarity gives us $\frac{L_p}{L_p}$ L_{η}
- Dynamic similarity in this case tells us that we must have $\frac{U_p^2}{\mu^2}$ U_m^2 L_p L_{η}
- So $\frac{v_p}{u}$ \boldsymbol{v}_{n}

- Now thinking about Re, $Re = \frac{\rho UL}{L}$ μ U_{\perp} $\mathcal V$
- We want: $Re_p = Re_m$
- le $\frac{U_m L_m}{N}$ v_n U_pL_p $\nu_{\bm p}$

• So we need
$$
\frac{U_m L_m}{\nu_m} = \frac{U_p L_p}{\nu_p}
$$

• Giving
$$
\frac{v_m}{v_p} = \frac{U_m}{U_p} \frac{L_m}{L_p} = \frac{1}{50\sqrt{50}} = 0.00283
$$

• So
$$
v_m = 2.83 \times 10^{-9} m^2 / s = 0.00283
$$
 cSt

- The closest liquid is mercury with $v=1.18\times10^{-7}$ m² /s (0.118 cSt).
- However mercury is not a suitable experimental fluid for this problem and we need to find another scale for the model.

 $1 cSt = 10^{-6} m²/s$

